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## Three-Nucleon Forces in Neutron Rich Isotopes (from O to Ni)

Collaborators:
A. Cipollone, CB, P. Navrátil:

Phys. Rev. Lett. 111, 062501 (2013)
V. Somà, A. Cipollone, CB, P. Navrátil, T. Duguet: Phys. Rev. C 89, $061301 R$ (2014)



## Modern realistic nuclear forces

Chiral EFT for nuclear forces:

|  | 2 N forces | 3 N forces | 4 N forces |
| :---: | :---: | :---: | :---: |
| $\mathrm{LO} \mathcal{O}\left(\frac{Q^{0}}{\Lambda^{0}}\right)$ |  |  |  |
| $\mathrm{NLO} \mathcal{O}\left(\frac{Q^{2}}{\Lambda^{2}}\right)$ |  | —— |  |

Single particle spectrum at $E_{\text {fermi }}$ :

[T. Otsuka et al. Phys Rev. Lett 105, 032501 (2010)]

Need at LEAST 3NF!!! ("cannot" do RNB physics without...)
$\mathrm{N}^{2} \operatorname{LO} \mathcal{O}\left(\frac{Q^{3}}{\Lambda^{3}}\right)$
(3NFs arise naturally at N2LO)

Saturation of nuclear matter:
 Phy.s Rev. C 88, 044302 (2013)]

## Chiral Nuclear forces - SRG evolved



## Faddeev-RPA in two words...

Particle vibration coupling is the main cause driving the distribution of particle strength-a least close to the Fermi surface...

these modes are all resummed exactly and to all orders in a ab-initio many-body expansion.

## Green's functions in many-body theory

One-body Green's function (or propagator) describes the motion of quasiparticles and holes:

$$
g_{\alpha \beta}(E)=\sum_{n} \frac{\left\langle\Psi_{0}^{A}\right| c_{\alpha}\left|\Psi_{n}^{A+1}\right\rangle\left\langle\Psi_{n}^{A+1}\right| c_{\beta}^{\dagger}\left|\Psi_{0}^{A}\right\rangle}{E-\left(E_{n}^{A+1}-E_{0}^{A}\right)+i \eta}+\sum_{k} \frac{\left\langle\Psi_{0}^{A}\right| c_{\beta}^{\dagger}\left|\Psi_{k}^{A-1}\right\rangle\left\langle\Psi_{k}^{A-1}\right| c_{\alpha}\left|\Psi_{0}^{A}\right\rangle}{E-\left(E_{0}^{A}-E_{k}^{A-1}\right)-i \eta}
$$

...this contains all the structure information probed by nucleon transfer (spectral function):


## Results for the N-O-F chains

A. Cipollone, CB, P. Navrátil, Phys. Rev. Lett. 111, 062501 (2013)

$\rightarrow$ 3NF crucial for reproducing binding energies and driplines around oxygen
$\rightarrow$ cf. microscopic shell model [Otsuka et al, PRL105, 032501 (2010).]

## Single nucleon transfer in the oxygen chain

[F. Flavigny et al, PRL110, 122503 (2013)]

## $\rightarrow$ Analysis of ${ }^{14} \mathrm{O}(\mathrm{d}, \mathrm{t})^{13} \mathrm{O}$ and ${ }^{14} \mathrm{O}\left(\mathrm{d},{ }^{3} \mathrm{He}\right)^{13} \mathrm{~N}$ transfer reactions @ SPIRAL





- Overlap functions and strengths from GF
- Rs independent of asymmetry


## Calcium isotopic chain

Ab-initio calculation of the whole Ca : induced and full3NF investigated


$\rightarrow$ induced and full3NF investigated
$\rightarrow$ genuine (N2LO) 3NF needed to reproduce the energy curvature and $\mathrm{S}_{2 n}$
$\rightarrow \mathrm{N}=20$ and $\mathrm{Z}=20$ gaps overestimated!
$\rightarrow$ Full 3NF give a correct trend but over bind!

## Neighbouring Ar, K, Ca, Sc, and Ti chains

V. Somà, CB et al. Phys. Rev. C89, 061301R (2014)

Two-neutron separation energies predicted by chiral NN+3NF forces:

$\rightarrow$ First ab-initio calculation over a contiguous portion of the nuclear chart-open shells are now possible through the Gorkov-GF formalism
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SUNERSEY

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Two-neutron separation energies predicted by chiral NN+3NF forces:


Lack of deformation due to quenched cross-shell quadrupole excitations
$\rightarrow$ First ab-initio calculation over a contiguous portion of the nuclear chart-open shells are now possible through the Gorkov-GF formalism
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SURREY

## Inversion of $d_{3 / 2}-s_{1 / 2}$ at $N=28$



FIG. 1. (color online) Experimental energies for $1 / 2^{+}$and $3 / 2^{+}$states in odd-K isotopes. Inversion of the nuclear spin is obtained in ${ }^{47,49} \mathrm{~K}$ and reinversion back in ${ }^{51} \mathrm{~K}$. Results are J. Papuga, et al., PRL 110, 172503 (2013)

## ${ }^{\text {AK }}$ K isotopes

Laser spectroscopy @ ISOLDE

Change in separation described by chiral NN+3NF:


ESPE: "centroid" energies

(Gorkov calculations at $2^{\text {nd }}$ order)

## Two-neutron separation energies for meutron rich $K$ isotopes


V. Somà, CB et al., in prep.

## Ca and Ni isotopic chains

Calculations based on ramps D, N1, N2:


$\rightarrow$ Large J in free space SRG matter (must pay attention to its convergence)
$\rightarrow$ Overall conclusions regarding over binding and $S_{2 n}$ remain but details change

IM-SRG results from H. Hergert
SUNERSEY

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## Ca and Ni isotopic chains

Difference of calculated BEs to the experiment for different masses:

$\rightarrow$ In general over binding per nucleon ( $E / A$ ) appear to stabilize above A~40-50 but more investigations are required.

## Conclusions

- What to did we learn about realistic chiral forces from ab-initio calculation ?
$\rightarrow$ Leading order 3NF are crucial to predict many important features that are observed experimentally (drip lines, saturation, orbit evolution, etc...)
$\rightarrow$ Experimental binding is predicted accurately up to the lower sd shell (A~30) but deteriorates for medium mass isotopes (Ca and above) with roughly $1 \mathrm{MeV} / \mathrm{A}$ over binding.
$\rightarrow$ This hints to the need of more repulsion in future generations of chiral realistic forces.



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## Thank you for your attention!!!!



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A. Cipollone, CB, P. Navrátil, Phys. Rev. Lett. 111, 062501 (2013)


$\rightarrow$ d3/2 raised by genuine 3NF
$\rightarrow$ systematic underestimation of radii

N3LOr(Assr500Mev/c) chiral NN interaction evolved to $2 \mathrm{~N}+3 \mathrm{~N}$ forces ( $2.0 \mathrm{fm}^{-1}$ )
N2SE(RR4OOMev/c) chiral 3N interaction evolved $\left(2.0 f \mathrm{~m}^{-1}\right)$

## Ca spectral distributions - at 2nid order

neu $\lambda_{\text {SRG }}=2 \Lambda=13,16,16$ RAMP Dpp FULL


## Pairing gaps

© Three-point mass differences

$$
\Delta_{n}^{(3)}(A)=\frac{(-1)^{A}}{2}\left[E_{0}^{A+1}-2 E_{0}^{A}+E_{0}^{A-1}\right]
$$



## Pairing gaps

Inversion of odd-even staggering

$\xrightarrow{\prime} \rightarrow$ Second order and 3NF necessary to invert the staggering

## Ni spectral alistributions - at 2tid order



## Collaborators

## SUNRRSEY

$\qquad$ cea


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